High-Resolution Electron Microscopy of High-Temperature Superconductors

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ABSTRACT We review recent results obtained at Northwestern using high-resolution electron microscopy to study high-temperature superconductors. While in general these materials form large, very perfect single crystal grains which display very few imperfections, there is also evidence of slip defects, amorphous regions, and order-disorder transformations. We also report that the gadolinium-based superconductors and in one case yttrium-based superconductors show evidence for some copper solid solubility in the form of copper-rich planar defects. The structure of a metastable trigonal polytype is also reported, as are the effects of electron beam and water vapor damage to the materials.

INTRODUCTION

The recent wave of high-temperature ceramic superconductors, most notably the triple perovskite YBa2Cu3O7 (generally abbreviated to 123) and its derivatives where the yttrium is replaced by almost any rare earth ion, have raised a number of important questions for electron microscopists. Perhaps of primary importance is why some preparations yield "good" superconductors, while others do not, and why bulk preparation techniques have to date yielded specimens that fail to match the current-carrying capacity of thin film specimens. Is this due to internal defects, secondary phases, or perhaps grain boundary phases in some of the materials? In addition, do these defects or secondary phases, if present, change as a function of processing of the ceramics?

For the last few months we have been fairly intensively studying these conventionally produced powder superconductors at Northwestern (Hwu et al., 1987a,b), using as our primary tool high-resolution electron microscopy supplemented with some microanalysis and diffraction. To a certain extent many of our results are somewhat negative; we have been unable to find any clear-cut structural features that could be responsible for poor behavior. What we have found are some interesting structural variations and defects which might in some cases lead to degradation of the properties. In this paper we will briefly review our results to date on the perfect superconductor, planar slip defects, copper-rich defects, order-disorder phenomena, polymorphs, and the effects of electron beam and water vapor damage. A short report on some of these results has already appeared (Buckett et al., 1987; Marks et al., 1987), and more detailed reports of other parts of the work are in preparation (Li and Zhang; Zhang et al.; Shibahara et al.)

EXPERIMENTAL METHOD

Superconductor specimens were prepared by the standard repeated annealing, grinding procedure and then cooled slowly in air (see, for instance, Hwu et al., 1987a,b). We have explored only the YBa2Cu3O7 compound and the related GdBa2Cu3O7 material which has been shown by a number of authors (e.g., Fisk et al., 1987; Hor et al., 1987; Hulliger and Ott, 1987) to be an equivalent superconductor. As a rule Meissner and conductivity measurements were used to check the properties of the specimens. To explore metastable and defective phases, superconducting specimens were prepared by quenching both in liquid nitrogen and oil. The specimens were then very gently ground (typ-
Fig. 1. Images of the triple perovskite structure taken along the three major zones, from top to bottom [100], [110], and [001], with diffraction patterns on the left.

ically between two glass plates) in order to minimize mechanical deformation, mounted on pyrolyzed carbon films, baked briefly on a light bulb for 10 minutes to reduce surface hydrocarbons, and then examined in either a Hitachi 700 at 200 kV or a Hitachi H-9000 operated at 300 kV. To check on the image interpretation, images were stimulated on
an Apollo workstation ring using programs written by one of us (L.D.M.) which are interface to SEMPER routines operating on the 1,024 x 1,024-pixel displays of Apollo 660 workstations. Conditions for the calculation were a slice thickness of 0.28 to 0.2 nm sampled to about 6 nm⁻¹ in reciprocal space with a Cs of 0.9 mm at 300 kV, a focal spread of 8 nm, and a convergence of 0.5 to 1.0 mr.

**Perfect superconductor**

The main component of the specimens prepared by slow cooling in oxygen is the perfect triple perovskite—see, for instance, Figure 1 and Marks et al. (1987). Typically traces of one or two impurity phases are present in a well-prepared specimen owing to slight deviations from the true 123 metal composition and/or incomplete mixing; less careful preparation in our experience introduced more impurity phases. Image simulations indicate that the material has the structure originally determined by X-ray (Hazen et al., 1987; Siegrist et al., 1987) and neutron diffraction (Beech et al., 1987; Beno et al., 1987; Greedan et al., 1987); the basic structure of the 123 compound is a solved question, so it is overkill to go to the length of very detailed image simulations to detect all the atomic positions. As observed by numerous authors (see, for instance, other papers in this volume and the papers by Beyers et al., 1987; Hewat et al., 1987; Hyde et al., 1987) the orthorhombic phase is normally twinned on (110) planes, presumably a stress-relief process during the tetragonal to orthorhombic phase transition; this can clearly be seen by splitting of the higher-angle diffraction spots. One exception to the general result of the perfect perovskite structure should be noted here; in thin film specimens (courtesy of workers at Stanford University) preliminary results indicate that there are extensive amorphous regions.

**Tetragonal phase**

As a rule, most of the specimen has the perfect structure shown above, and further analysis is unwarranted, as mentioned above, since X-ray and neutron diffraction methods are superior to imaging methods for perfect

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**Fig. 2.** Image taken along the [100] zone of the tetragonal phase.
structures. The thrust of our research, therefore, turned towards defects, either native to the structure or produced by varying the preparation conditions. The first experiment that was tried was liquid nitrogen quenching, which is known to produce a non-superconducting tetragonal structure with an average oxygen content of <6.5 atoms. The resultant structure (Fig. 2) is essentially the same as that of the orthorhombic phase, as also observed by X-ray methods (e.g., Gallagher et al., 1987).

**Slip defects**

In general, the material as mentioned above is rather perfect. One exception to this is that planar defects and in some cases cracks can often be found along (001) planes near to the surface of very thin flakes as shown in Figure 3. These defects can be understood as due to slip normal to the c-axis, caused by mechanical deformation during grinding to prepare the high-resolution specimens; only the thin flakes, not the thicker regions, show these defects. An earlier observation of similar defects by Ourmazd et al. (1987), who suggested that they could be the source of the superconductivity, can be attributed to use of room-temperature ion-beam milling to produce samples; we (Marks et al., 1987) along with a number of other authors (Hewat et al., 1987; Hyde et al., 1987) are confident that these defects are just an artifact. As further confirmation that the (001) plane is a preferred slip plane it should be noted that sinter-forged specimens are highly textured with the c-axis parallel to the pressing direction (Robinson et al., 1987).
Copper-rich defects

While the slip defects were the only defects that are common in the yttrium superconductor, two other defect structures can occur, both of which indicate that it is possible to obtain some copper solid solubility in the system. In one, and only one grain of liquid-nitrogen-quenched yttrium superconductor the defects shown in Figure 3a were observed. Quite commonly in the gadolinium superconductor we observed (Shibahara et al., in preparation) the defects shown in Figure 4b, 4c. The defects in Figure 4a,b correspond to an additional copper plane and a shear of approximately 1/6[331], essentially an antiphase domain boundary. The other defect shown in Figure 4c corresponds to a displacement along the c-axis with two Cu atoms in a line. At present there is no data on the GdBaCu oxide phase diagram so we cannot clearly understand the significance of the latter results, except to note the difference relative to the yttrium superconductor where it is established that the 123 compound is a stable point phase (e.g., Wang et al., 1987). It should be mentioned that Zandbergen et al. (1987) have reported similar defects in ErBa2Cu3O7 taken using a JEOL 200CX, describing them as additional planes of Er, which does not agree with our conclusions.

Order-disorder and polytypism

One of our first observations of these materials was that there existed classical order-disorder transformations [1]. The perfect structure has the Ba and Y atoms ordered on the (100) planes of the perovskite parent cell, but some regions (see Fig. 5) showed loss of the triple periodicity, implying random distribution of the Y and Ba atoms. (Initially this region showed some partially ordered regions when viewed using a TV pickup system, but these degenerated rapidly in the
electron beam.) This occurred in some grains from which the perfectly ordered phase appeared to be growing, although there were some indications from EDX analysis that these regions might be yttrium deficient.

A more recent result (Li and Zhang, in preparation) confirms the idea that the Y and Ba atoms need not be ordered in the structure of $\text{YBa}_2\text{Cu}_3\text{O}_{6.5}$ produced by quenching in oil (oil is a much faster quencher than liquid nitrogen). The resultant material (see Fig. 6) is a combination of very small regions of two trigonal phases with lattice parameters of $a=0.536$ and $c=0.666$ nm, the first having $P3\text{ml}$ symmetry with the Y atoms stacked along (111) planes of the parent perovskite, while the second has $P3\text{ml}$ symmetry and the same arrangement of atoms except that there is interchange of the Y atoms and one of the Cu sites within the unit cell as shown in Table 1. We believe that at high temperature the entropically favored structure will be a fully disordered structure with the yttrium and barium completely disordered in a simple perovskite structure. During quenching, as an alternative to the true 123 ordering along $[001]$, some ordering can occur on (111) planes leading to these metastable trigonal phases. On slower cooling, the lower energy ordering along $[001]$ takes place.

Electron beam and water damage

Although many of the better yttrium superconducting specimens are quite beam resistant, some, primarily those with a lower oxygen content, damage in the beam (Zhang et al., in preparation). The initial consequence is the formation of an amorphous surface phase (see Fig. 7a), which grows in from the surface. This is followed by segregation of the barium to the surface in the form of barium oxide (see Fig. 7b). The polytypic phase shown in Figure 6 also damages, with no amorphous phase formation but instead a loss of chemical ordering and segregation to
TABLE 1. Positions of the metal atoms in the two trigonal polymorphs of YBa$_2$Cu$_3$O$_{6.5}$.¹

<table>
<thead>
<tr>
<th>Atoms</th>
<th>Sites</th>
<th>x</th>
<th>y</th>
<th>z</th>
<th>Occupancy</th>
</tr>
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<tr>
<td>P3m1</td>
<td>Y</td>
<td>1(a)</td>
<td>0</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td></td>
<td>Ba</td>
<td>2(d)</td>
<td>1/3</td>
<td>2/3</td>
<td>1/3</td>
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<tr>
<td></td>
<td>Cu</td>
<td>1(b)</td>
<td>0</td>
<td>0</td>
<td>1/2</td>
</tr>
<tr>
<td></td>
<td>Cu</td>
<td>2(d)</td>
<td>1/3</td>
<td>2/3</td>
<td>5/6</td>
</tr>
<tr>
<td></td>
<td>O</td>
<td>3(e)</td>
<td>0</td>
<td>1/2</td>
<td>0</td>
</tr>
<tr>
<td></td>
<td>O</td>
<td>6(i)</td>
<td>1/3</td>
<td>1/6</td>
<td>1/3</td>
</tr>
<tr>
<td>P3ml</td>
<td>Y</td>
<td>1(b)</td>
<td>1/3</td>
<td>2/3</td>
<td>5/6</td>
</tr>
<tr>
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<td>1(b)</td>
<td>1/3</td>
<td>2/3</td>
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</tr>
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<td>6(i)</td>
<td>1/3</td>
<td>1/6</td>
<td>1/3</td>
</tr>
</tbody>
</table>

¹These data are based upon matching the diffraction patterns from single grains taken over a large range of angles and matching to the experimental images. The oxygen content was determined by RGA, and the oxygen positions have been taken from the basic perovskite lattice.
produce barium oxide on the surface (see Fig. 8). Two processes appear to be important, namely, diffusion of the barium atoms and oxygen DIET (desorption induced by electronic transitions; see, for instance, Petford et al., 1986). We suspect that some oxygen loss destabilizes the triple perovskite structure, which makes more space available for barium diffusion. Two other reports in the literature confirm this idea; Eaglesham et al. (1987) report an orthorhombic to tetragonal phase transition in the beam, which indicates oxygen loss, while Hyde et al. (1987) report the formation planar defects which would be expected from barium diffusion. Presumably, in the less oxygen-deficient samples the oxygen diffusion rate is slower, which reduces sharply the rate of damage. (Oxygen ejection solely from the immediate surface will have little effect unless the oxygen can be replenished by diffusion.)

These results are somewhat similar to the effect of water vapor on the material (Yan et al., 1987; Zhang et al., in preparation), which produces surface barium-rich needles (presumably barium carbonate originally) as shown in Figure 9. Both Auger and microanalysis show the surface barium, with microanalysis indicating that the perovskite frame is retained in a bulk barium-deficient
compound. We suspect that here again both oxygen loss and barium diffusion are important.

DISCUSSION

Perhaps the biggest problem for electron microscopists at present in the area of high Tc superconductors is that essentially all the results are negative; there is a clear negative correlation between the observed superconductivity and any crystalline defects. The current thinking in the literature is that the number and distribution of oxygen point defects is critical and oxygen unfortunately is not readily observable in the microscope. Although one can clearly in principle determine the difference between, for instance, the [100] and [010] directions in an image, the [loo] and [OlO] directions in an image, Tc superconductors is that essentially all the defects of the oxygen point defects will be hard if not impossible under the microscope. Perhaps the biggest problem for electron microscopy is yet to come, i.e., its use in microscopists at present in the area of high Tc superconductors is that essentially all the correlation between the observed superconductivity and any crystalline defects. The current thinking in the literature is that the number and distribution of oxygen point defects is critical and oxygen unfortunately is not readily observable in the microscope. Although one can clearly in principle determine the difference between, for instance, the [100] and [010] directions in an image, determining the order/disorder characteristics of the oxygen point defects will be hard if not impossible under the microscope. Perhaps the most important role of electron microscopy is yet to come, i.e., its use in characterizing processed materials which may contain unwanted phases or characterizing metal/superconductor composites. At present, however, processing of these materials seems to be a very dark area, and we have to await further progress.

ACKNOWLEDGMENTS

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REFERENCES


