

Electron Precession for Routine Crystallography

C.S. Own and L.D. Marks

Dept. of Materials Science, Northwestern University
2220 Campus Dr., Cook 2036
Evanston, IL 60208

Funding: Fannie and John Hertz Foundation, DOE, & UOP LLC

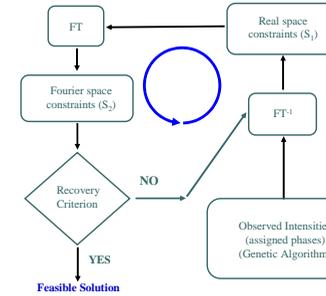
< Direct Methods (DM) >

In diffraction experiment we measure intensities (phase information is lost):

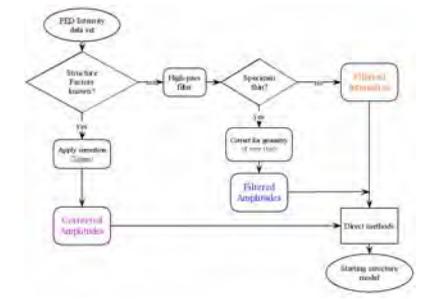
$$S.F. = F(k) \exp(-i\phi(k))$$

$$I(k) = |F(k)|^2$$

- If the phases were known:
 - Generate scattering potential maps
- For atomic structures, phase information is encoded within intensities
 - Requires < 10-20% RMS intensity error in order to work

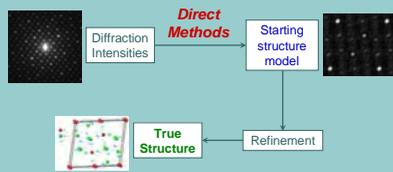


< PED Solution Flow Diagram >



< Introduction >

The difficult step in structure determination is obtaining a good starting structure



- If starting structure is close, refinement to true structure is accurate and reliable
- Direct methods** (phase recovery) is the key step but requires good intensities (see box)
- Traditionally, use X-rays (e.g., synchrotron):
 - Very **kinematical** (single scattering)
 - DM works well with kinematical data
- Transmission electron diffraction is better for many nanomaterials problems:
 - Fine probe (0.5Å-50nm)
 - Multiple data collection in one instrument
 - Chemical information (EELS, EDX)
 - Imaging and diffraction
 - In-situ* capabilities
 - TEM is widely available and can be less expensive

However, electron diffraction data is strongly distorted by **dynamical** (multiple) scattering:

- Intensity oscillates with specimen thickness
- Phase inversion on dynamical intensities is not usually successful
- A priori* information needed
 - Structure solutions to date using TEM usually need combination of multiple techniques
 - Time and \$\$\$**

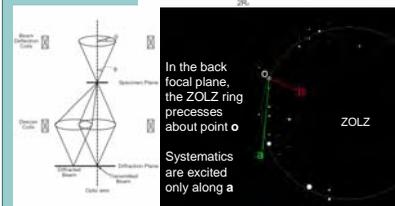
The solution:

- Precession Electron Diffraction (PED)**

< P.E.D.¹ >

- Conical beam rock/de-rock via deflectors
 - Avoids most multiple scattering – Usually only one strong beam excited at once.
 - An effective integration over excitation error
 - Like powder diffraction or texture patterns
 - Insensitive to tilt misorientation
 - Better intensity data via clever geometry**

Note: All results shown use parallel beam.



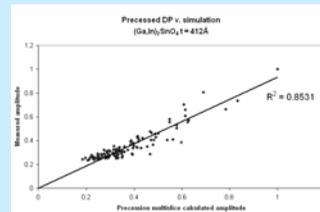
- Instrumentation¹⁰
 - PED mode retrofittable onto most any TEM
 - We have developed optical aberration compensations and rapid alignment⁹
 - The key variables^{8,10}:
 - Tilt angle ϕ
 - Thickness t
 - Spatial frequency of reflection $|g|$

Questions:

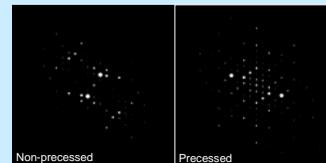
- Why were literature results inconsistent?⁴⁻⁶
- What is the true nature of the data?
- Do correction factors need to be applied, how and when should we apply them?

< Result 1: Precession Multislice >

I. Precession multislice is the correct model for precession:



Simulated patterns show improved intensity ordering.
($\phi = 24$ mrad, $t = 41$ nm, 200kV)

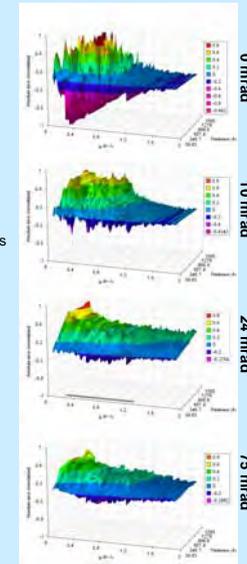


II. Amplitude Error plots



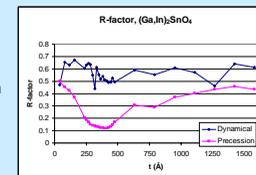
How to read these plots:

- ± 0.4 on plot corresponds to $\pm 16\%$ intensity error
- Bumpiness means data is dynamical
 - t-axis: pendellosung oscillations
 - g-axis: dynamical exchange between neighboring beams
- Trends as ϕ (prec. angle) \uparrow
 - Larger ϕ decreases error
 - Negative error disappears
 - After critical angle (~ 20 mrad), errors become small
 - Primary errors are near transmitted beam
 - Systematics, diffuse scattering
 - In regime $< 0.25 \text{ \AA}^{-1}$
 - Smaller thicknesses have low error
 - Will solve well with DM without correction factors
 - Larger thicknesses need correction



III. R-factor (for unrefined solution)¹¹

- Broad clear global minimum: $R_1 \sim 0.118$
 - Experiment matches simulated known structure
 - Compare to $R_1 \sim 0.3-0.45$ from previous precession studies after refinement (this is unrefined!)
- Accurate thickness determination:
 - Average $t \sim 41$ nm** (ordinarily intractable by DM)



$$R_1 = \frac{\sum (F_{exp} - F_{sim})}{\sum F_{exp}}$$

< Result 2: Correction Factors >

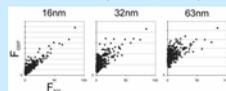
I. Previous corr. factor (an approx.)^{2,7}

Integral over all space (like powder diff.)

$$C_{blackman}(g, t, \phi) = g \cdot \left[1 - \sqrt{\frac{g}{2R_0}} \right] \times \left[\frac{A_g}{\int_0^t J_0(2x) dx} \right] \quad A_g = \frac{\pi}{g^2}$$

An *a priori* correction: Geometry-only (no dynamical)

Geometry term applied to calculated PED datasets (24 mrad, 200kV):



Result: Does not work! Need dynamical term (requires knowledge of structure factors)

II. New, more exact corr. factor

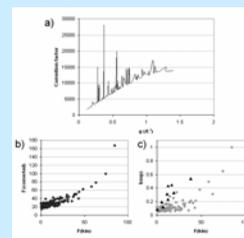
- Exact integral over scattered intensity
 - Sinc function altered by $\frac{g}{g_g}$ (2-beam)
 - A function of structure factor F_g
 - Some F_g must be known to use it!
 - $C_{blackman}$ converges to this for large t

$$C_{2beam}(g, t, \phi) = F_g^2 \left(\frac{1}{g} \int_0^t \frac{2g \sin^2(\frac{\pi s_{off}}{g})}{s_{off}^2} ds_{off} \right)^{-1}$$

$$s_{off} = \sqrt{s^2 + \frac{1}{g_g^2}} \quad \xi_g = \frac{\pi V_g \cos \theta_g}{\lambda F_g}$$

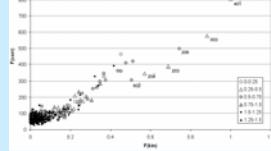
C_{2beam} applied to GITO system:

- Correction factor v. |g|
- Corrected multislice amplitudes
- Experimental intensities



Correction factor applied to simulated data
Plot of F_{corr} v. F_{kin} ($t = 127$ nm, $\phi = 75$ mrad)

(No apparent g-preference)



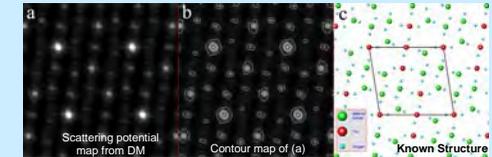
Result:

- An exact integration of two-beam approximates PED well
- Consider the Bessel integral limits
 - Precession intensity approximates true amplitude for large A_g

< Result 3: Examples >

I. (Ga,In)₂SnO₄ [010]³

- Using intensities (F_g^2) w/ DM
 - No filtering, all cation peaks found
- Displacement: $R_{neutron} - R_{precession}$
 - Average ΔR less than **4 pm!**
 - PED solution (unrefined) matches neutron-refined solution very closely¹¹

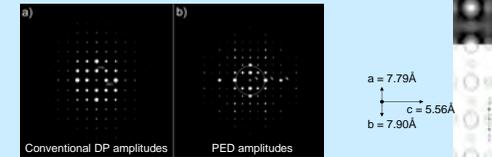


Atom position deviations (unrefined v. refined):

	ΔR (Å)
Sn1	0.00E+00
Sn2	0.00E+00
Sn3	6.55E-03
In/Ga1	5.17E-02
In/Ga2	2.37E-03
Ga1	6.85E-02
Ga2	1.22E-01

II. Al₂SiO₅ [110] (andalusite)

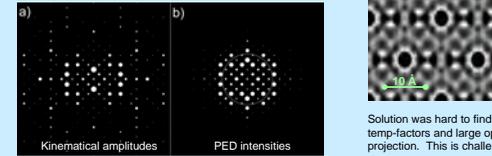
- Using amplitudes, high-pass filtering
 - Forbidden disappear in PED pattern
 - All cation peaks found, see vertical splitting (O features)



$a = 7.79 \text{ \AA}$
 $b = 7.90 \text{ \AA}$
 $c = 5.56 \text{ \AA}$

III. Al₈Si₄₀O₉₆ [001] (mordenite)

- Using intensities, high-pass filtering
 - Thick specimen, poor projection, damages easily



Solution was hard to find (above is a babinet soln.); large temp-factors and large open framework give poor projection. This is challenging for PED, but solvable.

< Conclusions / Future Work >

- Now we have better understanding of PED
 - Reduces overall error, errors @ low-index reflections
 - Precession extends usable thickness > 50 nm
 - Correction factor must include dynamical
 - 2-beam w/ exact integration limits is a sufficient model
 - Must exceed critical angle ($\phi > 20$ mrad)
- Structure maps from good PED expt.
 - DM maps with well-defined peaks
 - See cations, don't see light atoms
 - Can process much thicker specimens

Future Work

- More structures
 - Repertoire of solved structures
- Aberration corrected precession
 - Fine probe, fancy scanning configuration
 - Can avoid multi-beam excitation, also data mirroring
- A general correction factor (iterative)
 - For thick specimens

< References >

- Vincent & Midgely, Ultramic, 1994.
- Gjonnes, Ultramic, 1997.
- Sinkler, et al. J. Solid State Chem. 1998.
- Gemmi, et al. Acta Cryst. 1998.
- Gjonnes, et al. Acta Cryst 1998a.
- Gjonnes, et al. Acta Cryst 1998b.
- Blackman, Proc. Roy. Soc. 1999.
- Own, Subramanian, & Marks, Mic. & Microanal. 2004.
- Own & Marks, US Patent application # 60/531,641. Dec 2004.
- Own & Marks, Rev. Sci. Instr. 2005
- Own, Sinkler & Marks, Ultramic, in press.